

Metastability of the hard-core process on bipartite graphs

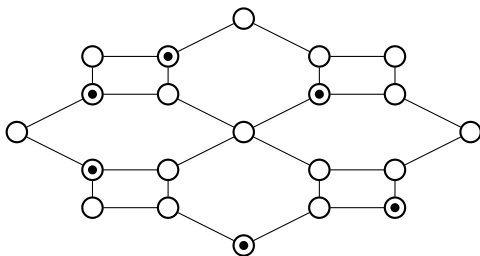
Frank den Hollander¹ Francesca Nardi² Siamak Taati¹

¹Mathematical Institute, Leiden University

²Department of Mathematics, Eindhoven University of Technology

METASTABILITY Workshop
Eurandom, April 2016

Hard-core gas process



Configurations

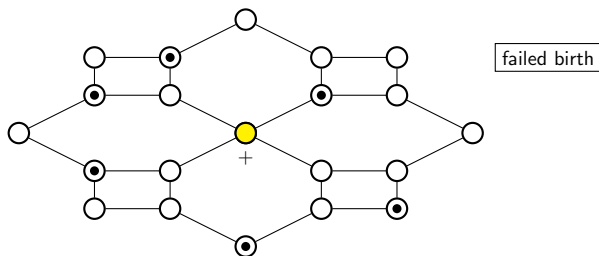
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- ▶ **Constraint:** particles cannot sit next to each other.

[Particles cannot overlap!]

Dynamics

- ▶ **Birth attempt** at site k (Poisson clock with rate λ_k)
- ▶ **Death attempt** at site k (Poisson clock with rate 1)
- ▶ All clocks are independent.

Hard-core gas process



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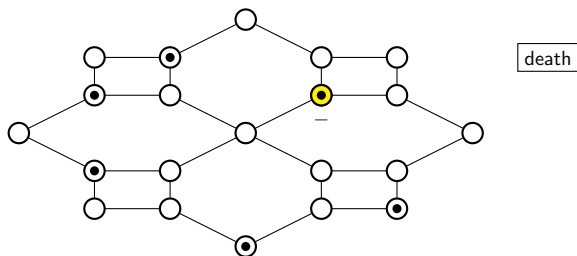
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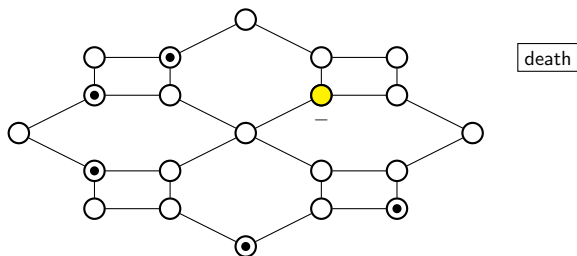
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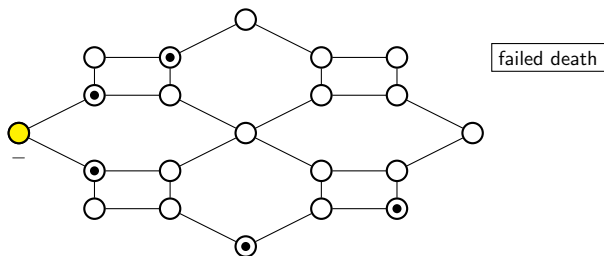
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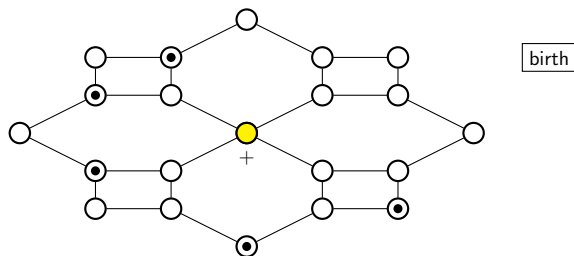
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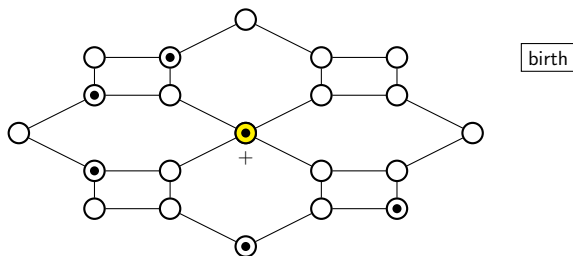
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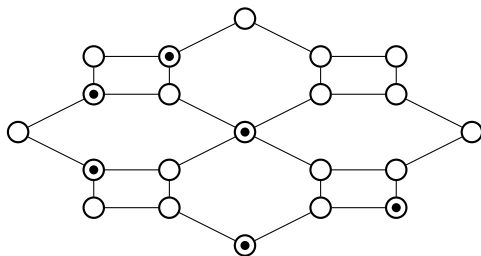
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Reversible stationary distribution

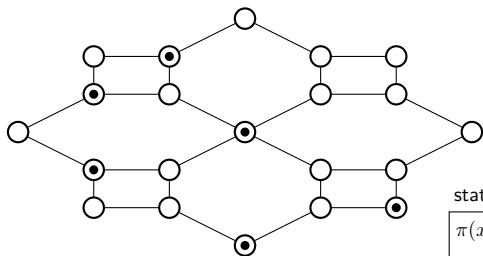
[Boltzmann distribution]

$$\pi(x) = \frac{1}{Z} \prod_{\substack{k \text{ occupied} \\ \text{in } x}} \lambda_k$$

for each **valid** configuration x .

(Z is the appropriate normalizing constant.)

Hard-core gas process



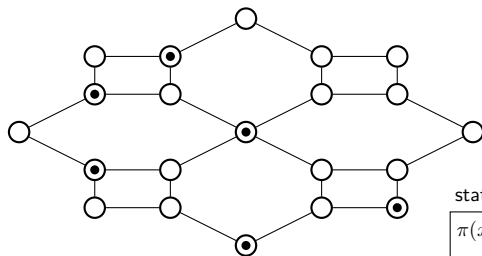
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Asymptotic regime

When the birth rates λ_k are very large:

Hard-core gas process



reversible
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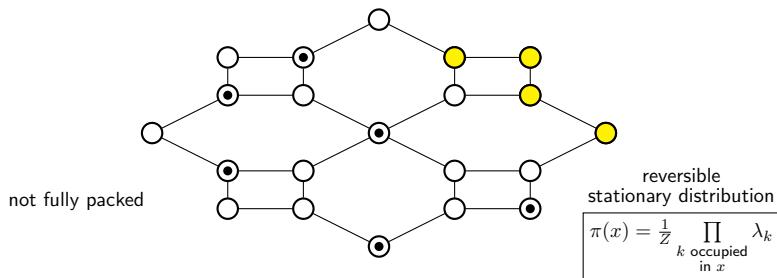
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When the **birth rates** λ_k are **very large**:

- ▶ The process tends to remain close to **fully packed** configurations, specially those that are "**locally stable**".

Hard-core gas process

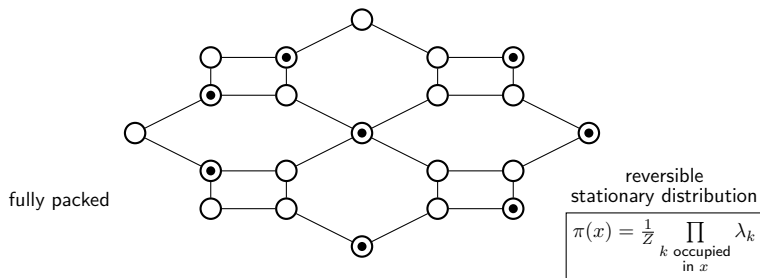


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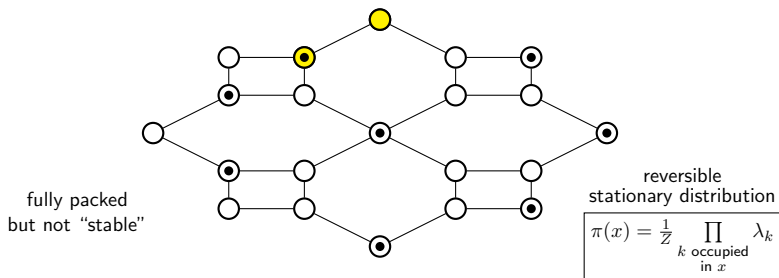


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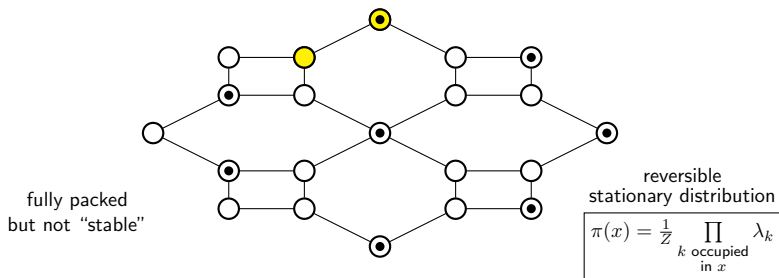


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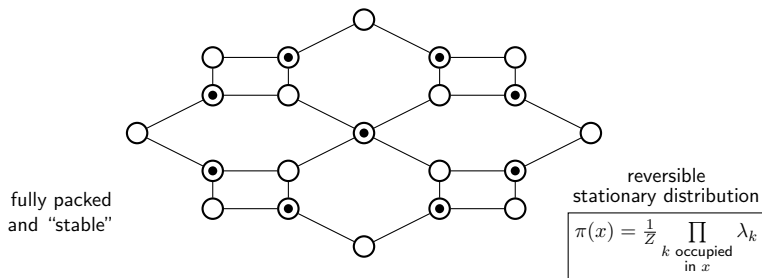


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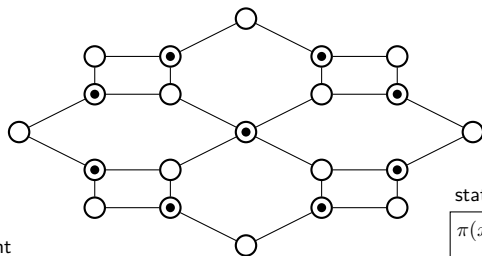


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Hard-core gas process



fully packed
and “stable”
but not efficient
(9)

reversible
stationary distribution

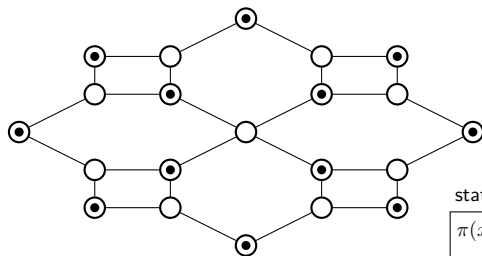
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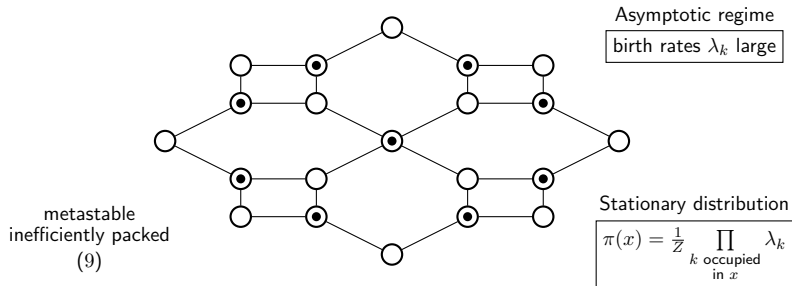
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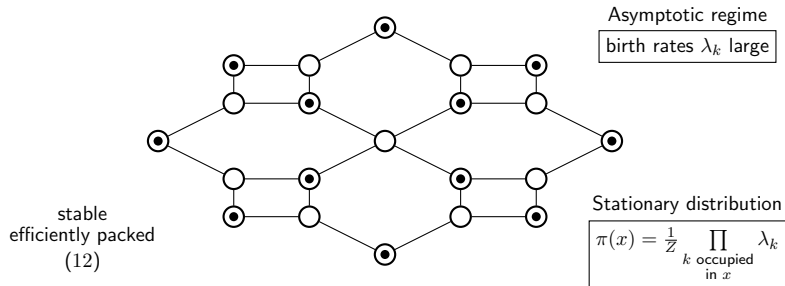
Hard-core gas process



Metastability

- ▶ It takes a long time for the process to leave a “**locally stable**” but **inefficiently packed** configuration. [large exit time]
- ▶ Once a more efficient configuration is reached, it takes much longer to return. [small stationary probability]

Hard-core gas process

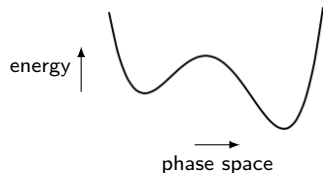


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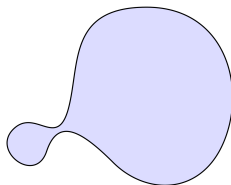
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Hard-core gas process

local minimum
in the energy landscape



bottleneck in
the phase space



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Hard-core gas on graphs

Motivation

- ▶ classic example from statistical mechanics [on the lattice]
 - phase transition (solid-gas) with symmetry breaking
- ▶ wireless communication networks
 - the graph represents the possibilities of interference
 - metastability undermines the network performance
- ▶ includes the Widom-Rowlinson model

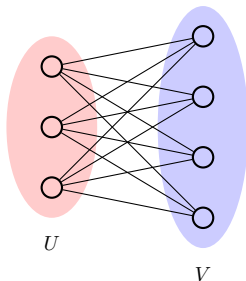
Related work

- ▶ Zocca, Borst, van Leeuwen and Nardi (2013–2015)
- ▶ Alessandro Zocca's PhD thesis (2015)
- ▶ Galvin and Tetali (2006), Randall (2008),
— and Antonio Blanca (2012)

Hard-core gas process

An exaggerated example

- ▶ complete bipartite graph
- ▶ birth rate λ at each site
- ▶ λ large
- ▶ $|U| < |V|$



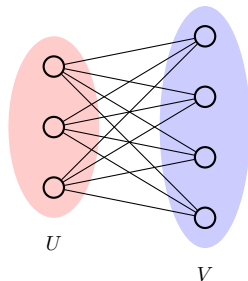
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- ▶ Exactly two fully packed configurations u and v
 - Both u and v are “locally stable”.
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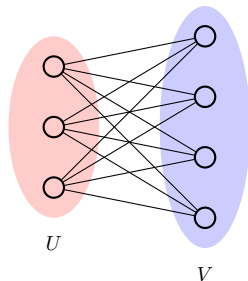
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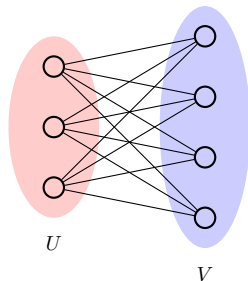
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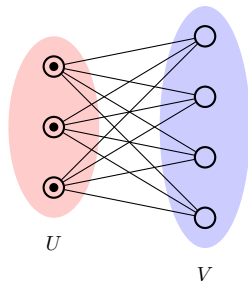
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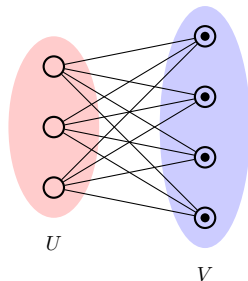
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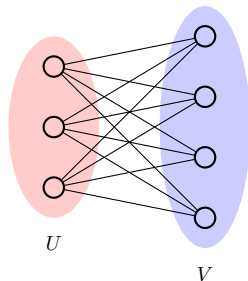
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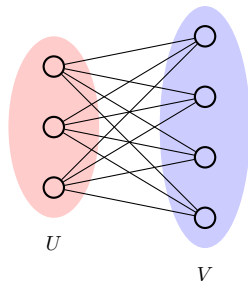
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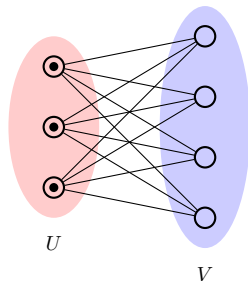
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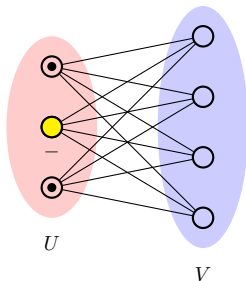
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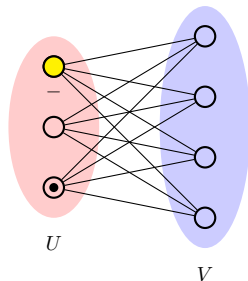
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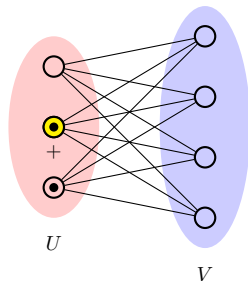
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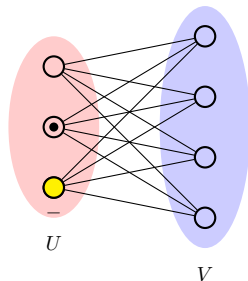
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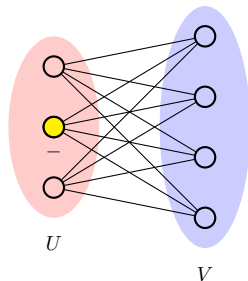
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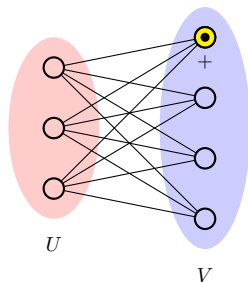
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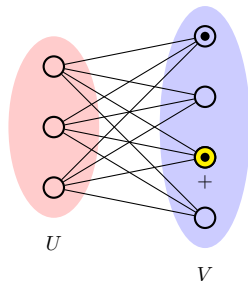
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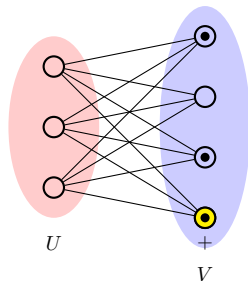
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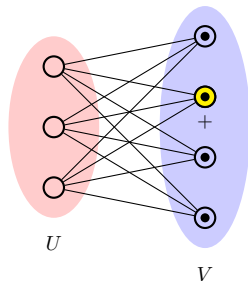
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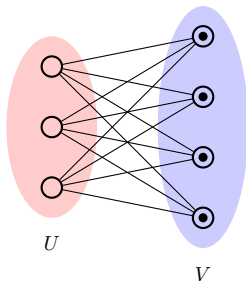
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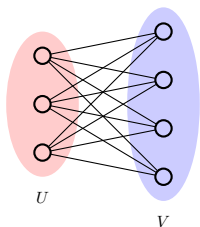
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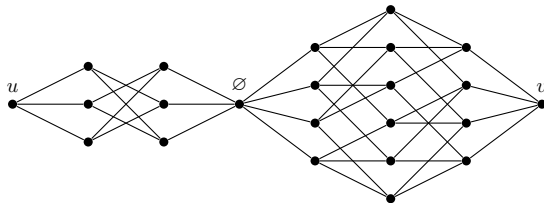
Question

How long does the transition from u to v take?

Hard-core process on a complete bipartite graph

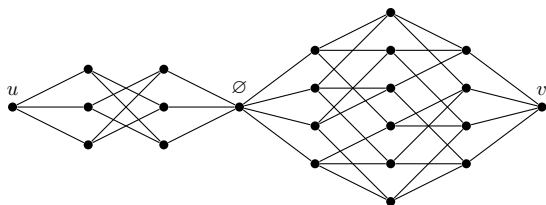


reversible Markov chain



Hard-core process on a complete bipartite graph

Reversible Markov chain

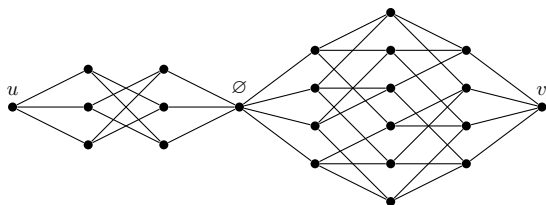


Let X_n ($n \geq 0$) be the **discrete-time** Markov chain.
The first hitting time of v is

$$T_v := \inf\{n \geq 0 : X_n = v\}.$$

Hard-core process on a complete bipartite graph

Reversible Markov chain

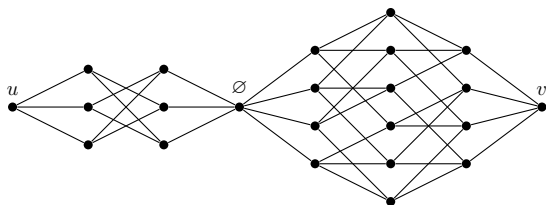


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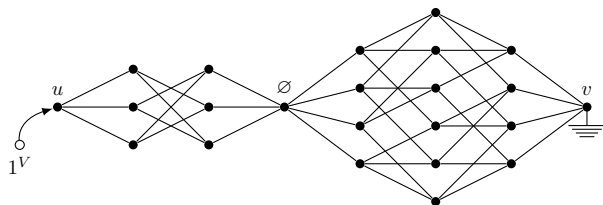
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Question

What is the *expected transition time* $\mathbb{E}_u T_v$?

Hard-core process on a complete bipartite graph

As an electric network



Review: reversible Markov chain vs. electric network

Fundamental connection I

For every state x ,

$$\mathbb{P}_x(T_A < T_B) = \text{voltage}(x)$$

if a 1^V battery is connected between A and B .

Fundamental connection II

For every state x ,

$$G_{T_B}(a, x) = \overbrace{\mathcal{R}(a \leftrightarrow B)}^{\text{effective resistance}} \pi(x) \mathbb{P}_x(T_a < T_B)$$

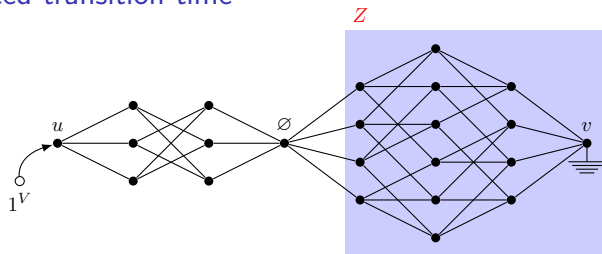
where $G_{T_B}(a, x) := \mathbb{E}_a[\# \text{ of visits to } x \text{ before } T_B]$.

Corollary

$$\mathbb{E}_a T_B = \mathcal{R}(a \leftrightarrow B) \sum_x \pi(x) \mathbb{P}_x(T_a < T_B)$$

Hard-core process on a complete bipartite graph

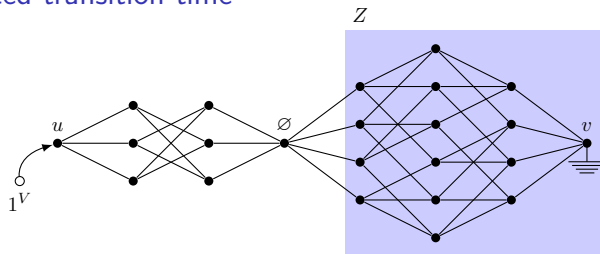
Expected transition time



$$\begin{aligned}\mathbb{E}_u T_v &\approx \mathbb{E}_u T_Z \\ &= \mathcal{R}(u \leftrightarrow Z) \sum_x \pi(x) \mathbb{P}_x(T_u < T_Z) \\ &= \pi(u) \mathcal{R}(u \leftrightarrow Z) \sum_x \frac{\pi(x)}{\pi(u)} \mathbb{P}_x(T_u < T_Z) \\ &\approx \pi(u) \mathcal{R}(u \leftrightarrow Z)\end{aligned}$$

Hard-core process on a complete bipartite graph

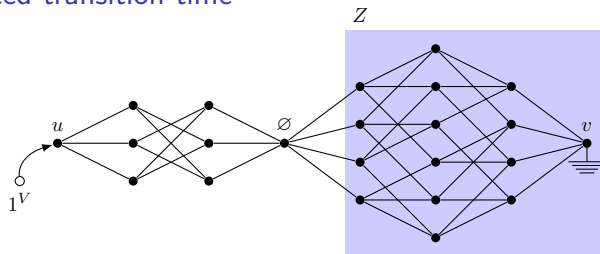
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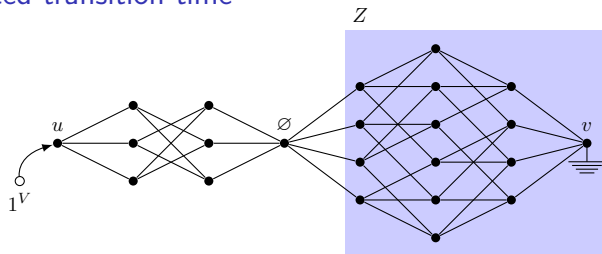
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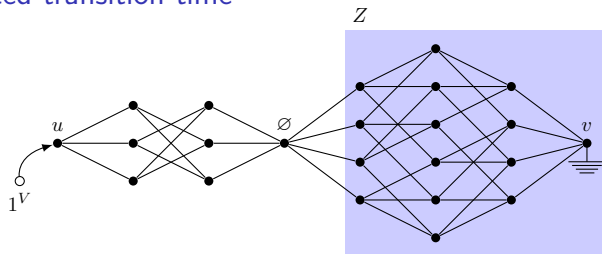
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Hard-core process on a complete bipartite graph

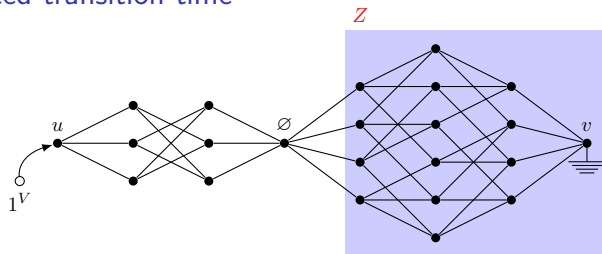
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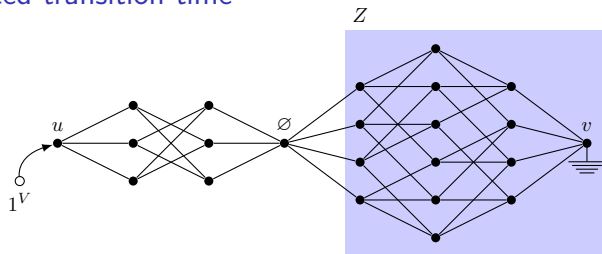
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$$\mathbb{E}_u T_v \approx \pi(u) \mathcal{R}(u \leftrightarrow Z) \approx \pi(u) \mathcal{R}(u \leftrightarrow v)$$

Hard-core process on a complete bipartite graph

Expected transition time

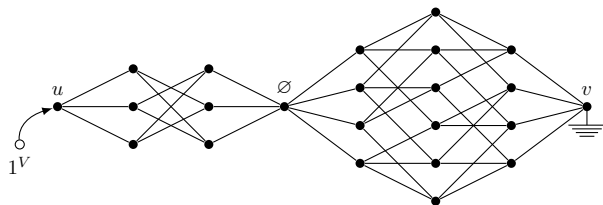


$$\mathbb{E}_u T_v \approx \pi(u) \mathcal{R}(u \leftrightarrow Z) \approx \pi(u) \mathcal{R}(u \leftrightarrow v)$$

It remains to estimate $\mathcal{R}(u \leftrightarrow v)$.

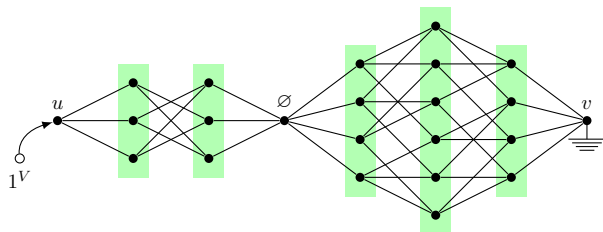
Hard-core process on a complete bipartite graph

Estimating the effective resistance



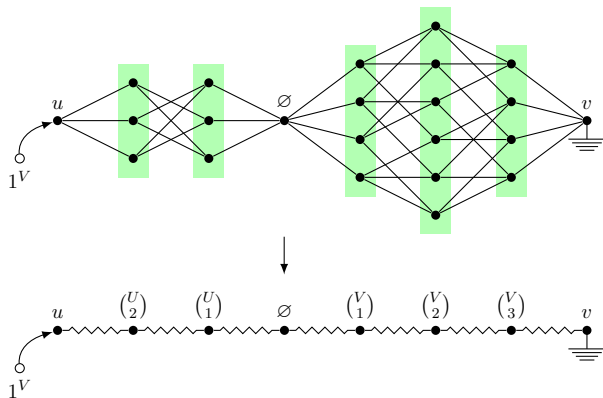
Hard-core process on a complete bipartite graph

Estimating the effective resistance



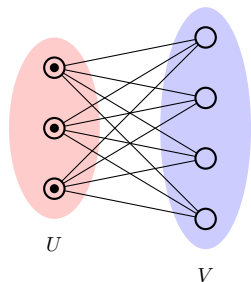
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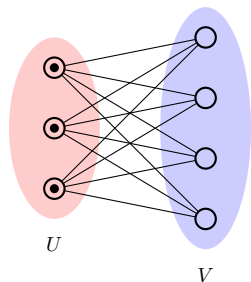


Proposition (Discrete time)

$$\mathbb{E}_u T_v = \frac{1}{|U|} \lambda^{|U|-1} [1 + o(1)] \quad \text{as } \lambda \rightarrow \infty.$$

Hard-core process on a complete bipartite graph

Expected transition time



Proposition (Continuous time)

$$\mathbb{E}_u T_v = \frac{\gamma}{|U|} \lambda^{|U|-1} [1 + o(1)] \quad \text{as } \lambda \rightarrow \infty.$$

$\gamma := (|U| + |V|)(1 + \lambda)$ is the rate of Poisson clock

Hard-core process on a bipartite graph

A more general setting

- ▶ an arbitrary bipartite graph

- ▶ Birth rates

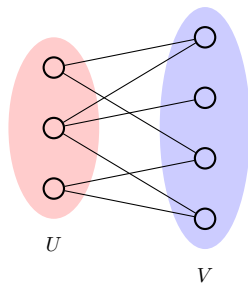
$$\lambda$$

on U

$$\bar{\lambda}$$

on V

- ▶ $\lambda, \bar{\lambda}$ large



Hard-core process on a bipartite graph

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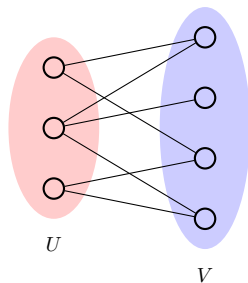
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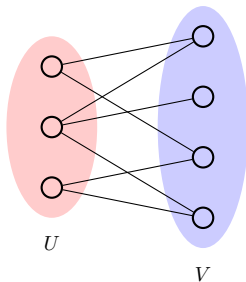
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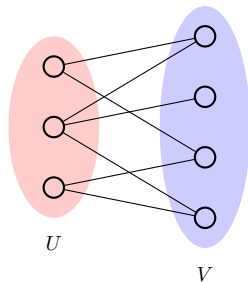
Hard-core process on a bipartite graph

A more general setting

- ▶ an arbitrary bipartite graph
- ▶ Birth rates (with $0 < \alpha < 1$)

$$\begin{array}{ll} \lambda & \text{on } U \\ \bar{\lambda} = \lambda^{1+\alpha+o(1)} & \text{on } V \end{array}$$

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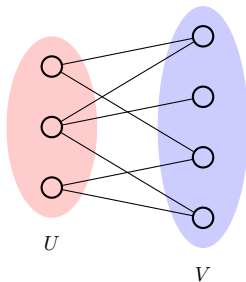
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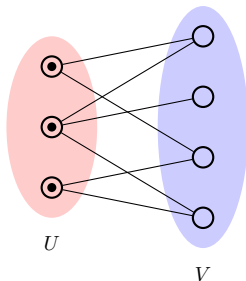
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Intuitive observations

- ▶ Two **fully packed** configurations u and v [but possibly many more]
 - Both u and v are “locally stable”.
 - v is the “most efficient” packing.

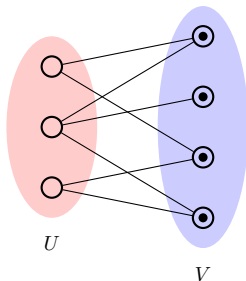
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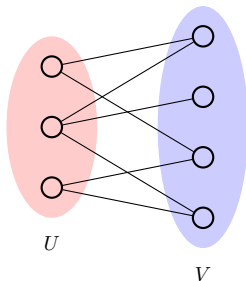
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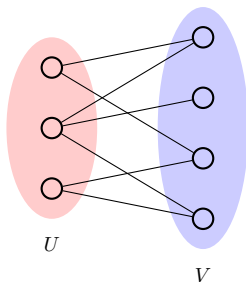
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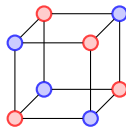
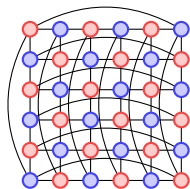
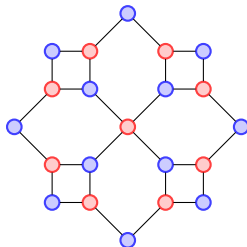
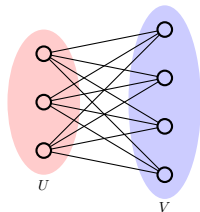


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Hard-core process on a bipartite graph

Examples of bipartite graphs



graphs arising from
two-species
Widom-Rowlinson model

Metastability in Markov processes

Some references

- ▶ Kramers (1940)
- ▶ **large deviations / path-wise approach**
 - ▷ Freidlin and Wentzell (1960–1970)
 - ▷ Cassandro, Galves, Olivieri and Vares (1984–)
 - ▷ ...
- ▶ **potential-theoretic approach**
 - ▷ Bovier, Eckhoff, Gaynard and Klein (2001–)
 - ▷ ...

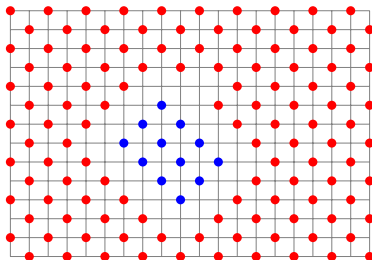
Three books

- ▶ Freidlin and Wentzell:
Random Perturbations of Dynamical Systems (1984)
- ▶ Olivieri and Vares: *Large Deviations and Metastability* (2005)
- ▶ Bovier and den Hollander:
Metastability — A Potential-Theoretic Approach (2015)

Main results: I

Theorem (Critical droplets)

For the hard-core dynamics on an even torus $\mathbb{Z}_m \times \mathbb{Z}_n$, when going *from u to v* , with large probability, the chain passes *through exactly one* transition $Q \rightarrow Q^*$, where Q and Q^* are obtained from the solutions of an *isoperimetric problem*.



A configuration in Q

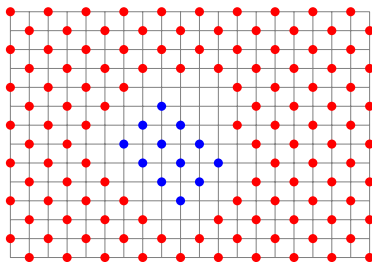
[similar for hypercube]

[similar for Widom-Rowlinson]

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A configuration in Q^*

[similar for hypercube]

[similar for Widom-Rowlinson]

Main results: II

Theorem (Expected transition time)

For the hard-core dynamics on an even torus $\mathbb{Z}_m \times \mathbb{Z}_n$ we have

$$\mathbb{E}_u T_v = \frac{\gamma}{2 m n l^*} \frac{\lambda^{l^*(l^*+1)+1}}{\bar{\lambda}^{l^*(l^*-1)}} [1 + o(1)]$$

as $\lambda \rightarrow \infty$, where $l^* := \lceil \frac{1}{\alpha} \rceil$ is the size of the critical droplet and $\gamma := |U|(1 + \lambda) + |V|(1 + \bar{\lambda})$ is the rate of the Poisson clock.

[similar for hypercube]

[similar for Widom-Rowlinson]

Proof steps.

Show that (in discrete time)

$$\mathbb{E}_u T_v = \pi(u) \mathcal{R}(u \leftrightarrow v) [1 + o(1)] \quad \text{as } \lambda \rightarrow \infty.$$

Estimate the effective resistance. □

Main results: III

Theorem (Asymptotic exponential law)

For the hard-core dynamics on “many” bipartite graphs we have

[e.g., torus, hypercube, ...]

$$\mathbb{P}_u \left(\frac{T_v}{\mathbb{E}_u T_v} > t \right) \rightarrow e^{-t}$$

uniformly in $t \in \mathbb{R}^+$ as $\lambda \rightarrow \infty$.

Intuition.

Many many trials (attempts to form a critical droplet) with tiny probability of success

\implies success time approximately exponential



Effective resistance: rough estimate

Critical resistance

[a.k.a. communication height]

For every two states $a, b \in \mathcal{X}$, set

$$\Psi(a, b) := \inf_{\omega: a \rightsquigarrow b} \sup_{e \in \omega} r(e)$$

Remark

- ▶ $a, b \mapsto \mathcal{R}(a \leftrightarrow b)$ is a *metric* on \mathcal{X} .
- ▶ $a, b \mapsto \Psi(a, b)$ is an *ultra-metric* on \mathcal{X} .

Proposition (Equivalence)

There exists a constant $k \geq 1$ such that

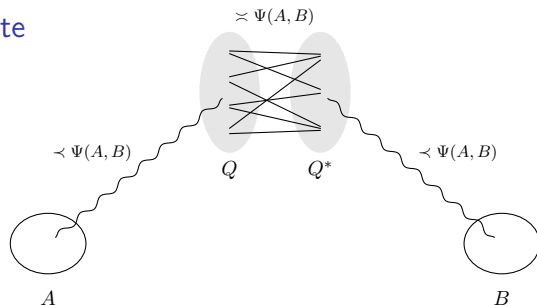
[independent of λ]

$$\frac{1}{k} \Psi(a, b) \leq \mathcal{R}(a \leftrightarrow b) \leq k \Psi(a, b)$$

for all $a, b \in \mathcal{X}$.

Effective resistance: sharp estimate

Critical gate

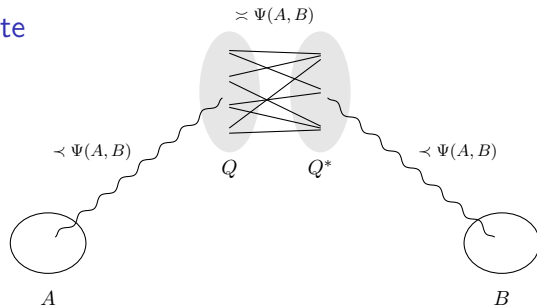


A pair (Q, Q^*) is a **critical gate** between A and B if

1. $r(x, y) \asymp \Psi(A, B)$ for every $x \in Q$ and $y \in Q^*$ with $x \sim y$,
2. $\Psi(A, x) \prec \Psi(A, B)$ for every $x \in Q$,
3. $\Psi(y, B) \prec \Psi(A, B)$ for every $y \in Q^*$, and
4. every **optimal path** from A to B passes through a transition $Q \rightarrow Q^*$.

Effective resistance: sharp estimate

Critical gate



Proposition

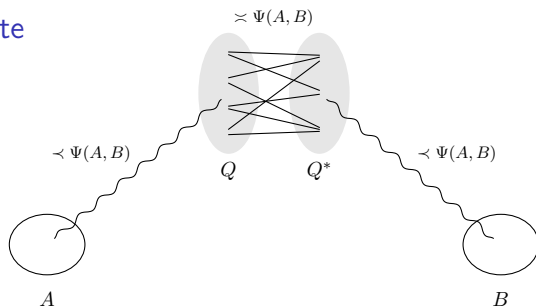
Let (Q, Q^*) be a critical pair between A and B . Then,

$$C(A \leftrightarrow B) = c(Q, Q^*) [1 + o(1)] \quad \text{as } \lambda \rightarrow \infty,$$

where $c(Q, Q^*) := \sum_{\substack{x \in Q \\ y \in Q^* \\ x \sim y}} c(x, y)$.

Effective resistance: sharp estimate

Critical gate



$$\mathcal{C}(A \leftrightarrow B) = c(Q, Q^*) [1 + o(1)]$$

Proof.

Upper bound: simple Nash-Williams inequality

Lower bound: generalized Nash-Williams inequality

[a.k.a. Berman-Konsowa variational principle]



Thank you for your attention!